



Analyse de machines tournantes comportant des non-linéarités avec Cast3M

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2016-10-10, GDR DYNOLIN

Plan

I. Motivation

II. Intégration temporelle

III. Equilibrage harmonique

IV. Conclusions et Perspectives

Contexte

- Etude et dimensionnement de machines tournantes
- Géométrie parfois complexe



Turboalternateur EDF

Développement d'outils dans logiciel EF

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Gas Turbine Modular Helium Reactor



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► Choix de Cast3M ⇒maîtrise et pérennisation



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- Choix de Cast3M ⇒maîtrise et pérennisation
- Analyses modales (modes \mathbb{R} et \mathbb{C}) et harmoniques (balourd ...)



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- Choix de Cast3M ⇒maîtrise et pérennisation
- ► Analyses modales (modes \mathbb{R} et \mathbb{C}) et harmoniques (balourd ...)
- Modélisations : poutre/2D Fourier/3D, massif/coque, isotrope/orthotrope ...



Présence de composants au caractère rapidement non-linéaire (paliers, squeeze-film ...)

⇒Analyses non-linéaires nécessaires !

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Intégration temporelle

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- Intégration temporelle
- Analyses dans le domaine fréquentiel (HBM)

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Equations of motion and Newmark scheme Newmark implicit Newmark explicit Taking into account the non-linearity

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Equations of motion

Rotor in rotating frame

 $[M]\underline{\ddot{u}}' + [\Omega G + C_{\textit{visc}}]\underline{\dot{u}}' + [K_{\textit{elas}} + \Omega^2 K_{\textit{cent}} + K(\sigma^0)]\underline{u}' = F'_{\textit{ext}}$

Linking rotating and non-rotating frame

$$\begin{pmatrix} U1\\U2\\U3 \end{pmatrix} = \begin{bmatrix} 1 & 0 & 0\\0 & \cos\Omega t & -\sin\Omega t\\0 & \sin\Omega t & \cos\Omega t \end{bmatrix} \cdot \begin{pmatrix} U1'\\U2'\\U3' \end{pmatrix}$$
$$U = R(\Omega t) \cdot U'$$

A Lagrange multiplier Λ is introduced $% \lambda$:

$$\begin{pmatrix} 0 & -I & R \\ -I^T & 0 & 0 \\ R^T & 0 & 0 \end{pmatrix} \cdot \begin{bmatrix} \Lambda \\ U \\ U' \end{bmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix}$$
$$L(\Omega t) \quad \cdot \quad q = 0$$



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The Newmark scheme $\ddot{u}_{n+1} = \frac{1}{\beta \Delta t^2} (u_{n+1} - u_n) - \frac{1}{\beta \Delta t} \dot{u}_n + (1 - \frac{1}{2\beta}) \ddot{u}_n$ $\dot{u}_{n+1} = \frac{\gamma}{\beta \Delta t} (u_{n+1} - u_n) + (1 - \frac{\gamma}{\beta}) \dot{u}_n + \Delta t (1 - \frac{\gamma}{2\beta}) \ddot{u}_n$

Schemes properties

			$\Delta t_c / \Delta t^e$
	1/2		
Fox and Goodwin	1/2	1/12	2.45
	1/2	1/6	
Average acceleration	1/2	1/4	00
Modified average acceleration	$1/2 + \alpha$	$1/4(1 + \alpha)^2$	∞
Runge-Kutta 4			
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Newmark average acceleration method (implicit)

Application to a linear 4-dof stator-rotor test problem



 $Kq + C\dot{q} + M\ddot{q} = F$

Kinematics constraint $L(\Omega t)q = 0$

avec : $q^T = \begin{pmatrix} \lambda_y & \lambda_z & u_y & u_z & u'_y & u'_z \end{pmatrix}$



Prabel, B. Some remarks on time integration of 3D rotor-stator assembly, ECCOMAS 2016

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$$q^{T} = (\lambda_{y} \quad \lambda_{z} \quad u_{y} \quad u_{z} \quad u'_{y} \quad u'_{z})$$

$$K = \begin{bmatrix} 0 & 0 & 1 & 0 & -\cos\Omega t & \sin\Omega t \\ 0 & 0 & 1 & -\sin\Omega t & -\cos\Omega t \\ k & 0 & 0 & 0 \\ k & 0 &$$



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 \Rightarrow instability of the "unconditionally stable scheme" !

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 \Rightarrow instability of the "unconditionally stable scheme" ! \Rightarrow Lagrange multipliers are suspected...

Newmark average acceleration method (implicit)

After long thinking...



Newmark average acceleration method (implicit)

... the solution was written !



Cardona, A., & Geradin, M. (1989). Time integration of the equations of motion in mechanism analysis. Computers & structures, 33(3), 801-820

"Newmark trapezoidal rule is unconditionally unstable in the presence of constraints" (analysis based on the computation of the eigenvalues and eigenvectors of the amplification matrix A of a constrained system)

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 \Rightarrow Proposed solution : use time integrator with controlled numerical damping.

Implicit schemes with numerical damping

Hilber-Hughes-Taylor and $\alpha\text{-generalized}$ methods

 $\begin{aligned} (1 - \alpha_m) M \ddot{u}_{n+1} + \alpha_m M \ddot{u}_n + (1 - \alpha_f) C \dot{u}_{n+1} + \alpha_f C \dot{u}_n + (1 - \alpha_f) K u_{n+1} + \alpha_f K u_n \\ &= (1 - \alpha_f) F_{n+1}^{ext} + \alpha_f F_n^{ext} \end{aligned}$
Implicit schemes with numerical damping

Hilber-Hughes-Taylor and α -generalized methods

$$(1 - \alpha_m)M\ddot{u}_{n+1} + \alpha_mM\ddot{u}_n + (1 - \alpha_f)C\dot{u}_{n+1} + \alpha_fC\dot{u}_n + (1 - \alpha_f)Ku_{n+1} + \alpha_fKu_n$$
$$= (1 - \alpha_f)F_{n+1}^{ext} + \alpha_fF_n^{ext}$$



Implicit schemes with numerical damping

Results of the 4-dof test problem



© No more instability

© Small dependance of the solution to the numerical damping

Implicit schemes with numerical damping

Results of the 4-dof test problem



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Implicit schemes with numerical damping

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Implicit schemes with numerical damping

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Implicit schemes with numerical damping

Application to a linear 3D rotor-stator assembly Mode shape constituting the CMS base :



Implicit schemes with numerical damping Application to a linear 3D rotor-stator assembly



No more instability
Dependance of the solution to the numerical damping

Implicit schemes with numerical damping





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Implicit schemes with numerical damping





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Newmark central difference method (explicit)

Simple 3D rotor test problem



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Cause of numerical instability ?

Newmark central difference method (explicit)

Cause of numerical instability ?

"Classical" simplification in explicit software

$$\left[M + \frac{\Delta t}{2}C\right]\ddot{u}_{n+1} = F_{n+1}^{ext} - Ku_{n+1} - C\dot{u}_{n+\frac{1}{2}}$$

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$$\begin{bmatrix} M + \frac{t}{2}C \\ \ddot{u}_{n+1} = F_{n+1}^{ext} - Ku_{n+1} - C\dot{u}_{n+\frac{1}{2}} \end{bmatrix}$$

M diagonal \Rightarrow no matrix inversion !

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Test case of Zhu et al.



Grandeur	а	b	С	d
$\alpha = \frac{m_{bearing}}{m_{disc}}$	0.1	0.267	0.1	0.1
$K = \frac{k_{spring}}{k_{shaft}}$	0.05	0.287	0.25	0.01
$\xi = \frac{c_{disc}}{2m_{disc}w_0}$	0.0005	0.0005	0.0005	0.0005
$U = \frac{F^{unbalance}}{\Omega^2 m_{disc}C}$	0.2	0.45	0.2	0.3
$B = \frac{\mu R L^3}{m_{bearing} w_0 C^3}$	0.05	0.145	0.01	0.025
and $w_0 = \sqrt{k_{shaft}/m_{disc}}$				

- Unbalance force
- Nonlinear force:

• Goal = determine the response for every speed of rotation Ω

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due to pressure distribution in squeeze-film (short bearing assumption):

$$p(\theta, z) = -\frac{6\mu C}{h^3} \left(\frac{L^2}{4} - z^2\right) \left(\dot{\epsilon}\cos\theta + \epsilon\dot{\phi}\sin\theta\right) \text{ with } h = C(1 + \epsilon\cos\theta)$$

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Newmark central difference method (explicit)

Référence (Zhu et al.)



Newmark central difference method (explicit)



Newmark central difference method (explicit)



- ▶ Implicit integration ⇒numerically damped scheme
- Explicit integration \Rightarrow "full" scheme
- ► It takes often a long (computational) time to reach steady state
- \Rightarrow Choice of efficient explicit computation on modal basis (DYNE operator of Cast3M) \odot
- ⇒ An alternative would be the computation of the nonlinear response in frequency domain...

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Intégration dans Cast3M Exemples académiques Flexible rotor on nonlinear squeeze-film damper Bifurcation tracking

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HBM

Intégration dans Cast3M

♦ Pb to solve: $-r = M\ddot{u} + C\dot{u} + Ku - f^{nl}(\dot{u}, u) - f^{ext} = 0$ ⇒discretization, model definition, boundary condition, loading, eigenmodes basis, ...

HBM

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♦ Fourier serie decomposition: $u(t) = U_0 + \sum_{k=1..H} U_k \cos k\omega t + V_k \sin k\omega t$ $-\mathbf{R}(\mathbf{U}, \omega) = Z(\omega)\mathbf{U} - \mathbf{F}^{n/}(\mathbf{U}) - \mathbf{F}^{ext} = 0$

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♦ Continuation solving method:



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♦ Continuation solving method:



 $\Rightarrow \mathsf{CONTINU} + \mathsf{AFT} \ \mathsf{procedur}$

♦ Stability analysis (Floquet exponents with Hill's method):

$$\left[\Delta_{0}+\lambda\Delta_{1}+\lambda^{2}\Delta_{2}\right]\tilde{\boldsymbol{U}}=\boldsymbol{0}$$

 \Rightarrow FLOQUET procedur

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Duffing oscillator

Duffing equation of motion:

$$\ddot{u} + 2\xi\omega_0\dot{u} + \omega_0^2u + \mu u^3 = p\omega_0^2\cos\omega t$$

Fixed values: $\xi = 5\%$, $\omega_0 = 1$ and p = 0.5. Continuation with respect to ω .

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ω/ω_

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Duffing oscillator

Duffing equation of motion:

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Forced Van der Pol oscillator

• Forced Van der Pol equation of motion: $\ddot{u} - \alpha (1 - u^2) \dot{u} + u = p \cos \omega t$

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Forced Van der Pol oscillator

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Forced Van der Pol oscillator

• Phase portrait for p = 0.5:



Forced Van der Pol oscillator

• Phase portrait for p = 0.5:



Plan

I. Motivation

II. Intégration temporelle

III. Equilibrage harmonique
Intégration dans Cast3M
Exemples académiques
Flexible rotor on nonlinear squeeze-film damper
Bifurcation tracking

IV. Conclusions et Perspectives

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0. Calcul de la courbe de réponse $\begin{pmatrix} \mathbf{U} & \boldsymbol{\omega} \end{pmatrix}^{T}$

1. Localisation des points de bifurcation $\begin{pmatrix} \mathbf{U} & \boldsymbol{\omega} & \boldsymbol{\phi} \end{pmatrix}^{T}$

2. Suivi des points de bifurcation

$$\begin{pmatrix} \mathbf{U} & \omega & \phi & \alpha \end{pmatrix}'$$

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Xie, L. Baguet S., Prabel, B. and Dufour, R., Numerical Tracking of Limit Points for Direct Parametric Analysis in Nonlinear Rotordynamics, Journal of Vibration and Acoustics, 138(2)-2016 Xie, L. Baguet S., Prabel, B. and Dufour, R., Bifurcation tracking by Harmonic Balance Method for performance tuning of nonlinear dynamical systems, Mechanical Systems and Signal Processing, accepted

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New unknown : parameter α Prediction : along ϕ New equation (correction steps) : pseudo-arc-length equation :

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Example: Jeffcott rotor (2 dof system) submitted to contact with friction



2. Suivi des points de bifurcation

Response for $\mu = 0.05$ Stable periodic Unstable periodique ----6 ΙP 5 Amplitude r/h 4 3 0.2 0.4 0.6 0.8 1.2 Frequency $\overline{\omega}$

2. Suivi des points de bifurcation

Response for $\mu = 0.11$



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2. Suivi des points de bifurcation

Response for $\mu = 0.20$



2. Suivi des points de bifurcation

Response for different values of $\boldsymbol{\mu}$



2. Suivi des points de bifurcation

Limit Point tracking (new parameter = μ)



2. Suivi des points de bifurcation

Neimark-Sacker tracking (new parameter = μ)



2. Suivi des points de bifurcation

LP and NS tracking


HBM

2. Suivi des points de bifurcation

LP and NS tracking



Frequency $\overline{\omega}$

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- [©] HHT and α -generalized implicit schemes implemented in Cast3M (only for linear analysis)
- Efficient explicit schemes available for nonlinear rotordynamics in Cast3M
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Future : Change \odot into \odot and explore collocation method, quasi-periodic motion computation, ...

Thank you for your attention !



Any question ?